

# Explorations of Holographic Kondo Physics: Quantum Quenches

Jackson Wu

University of Alabama

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**Based on works:**

J. Erdmenger, C. Hoyos, A. O'Bannon, JW  
**JHEP 1312, 086 (2013), arXiv:1310.3271**

J. Erdmenger, M. Flory, M. Newrzella, M. Strydom, JW  
**In progress**

# Advertisement

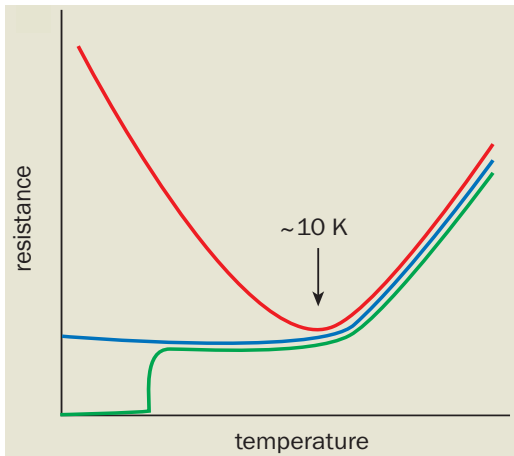
**“Quasinormal modes, charged magnetic brane in  $AdS_5$ , and (chiral) transport”**

M. Ammon, M. Kaminski, R. Koirala, J. Leiber, JW

# Outline

- Introduction
- The holographic Kondo model
- Quantum quenches
- Discussion

# The Kondo Effect



# The Kondo Effect

A **many-body** state (**Kondo resonance**) formed when conduction electrons hybridize with the (magnetic) impurity:

- Impurity screened
- **IR physics** changed

**Modern perspective:**

- RG flow triggered by a marginally relevant operator
- **Single-impurity problem:** rare example of flow between two trivial fixed points
- **Generalised Kondo problem:** **non-Fermi liquid IR FP!**

# Motivation

**Many techniques** developed to study Kondo physics:

- NRG, DMRG, Quantum Monte Carlo
- Bethe ansatz/Integrability
- Large-N
- CFT

**Non-equilibrium aspects** not easy to access with traditional techniques

**Holography** provides **new/alternative** method of study

# Holographic Description

## Essential ingredients of the Kondo problem:

- Chiral fermions:  $J^a = \psi_L^\dagger \Gamma^a \psi_L$
- Impurity:  $S^a = \chi^\dagger T^a \chi$  (slave fermion representation)
- Kondo interaction:  $J \cdot S = |\mathcal{O}|^2 + O(1/N)$ ,  $\mathcal{O} = \psi_L^\dagger \chi$



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## Single-impurity problem effectively 1D: CFT approach

- $J^a$  satisfy Kac-Moody algebra
- $\chi^\dagger \chi = Q \leftrightarrow R_{imp}$
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How to describe these holographically?

## Top-Down View

	$x^0$	$x^1$	$x^2$	$x^3$	$x^4$	$x^5$	$x^6$	$x^7$	$x^8$	$x^9$
$N_c$ D3	•	•	•	•	–	–	–	–	–	–
$N_7$ D7	•	•	–	–	•	•	•	•	•	•
$N_5$ D5	•	–	–	–	•	•	•	•	•	–

**Maldacena limit:**  $N_c \rightarrow \infty$ ,  $g_s \rightarrow 0$ ,  $\lambda \equiv 4\pi g_s N_c$  fixed;  $\lambda \rightarrow \infty$

**Probe limit:**  $\lambda_{5,7} \propto N_{5,7}/N_c \rightarrow 0$

- 3-7 strings: (1+1)-dim. chiral fermions  $\psi_L$
- 3-5 strings: (0+1)-dim. slave fermions  $\chi$
- 5-7 strings: tachyon

# Gravity Dual

## Near horizon limit:

- D3:  $AdS_5 \times S^5$  background
- D7 along  $AdS_3 \times S^5$ : Chern-Simons gauge field  $A$
- D5 along  $AdS_2 \times S^4$ : gauge field  $a$ , bifundamental scalar  $\Phi$

## Duality dictionary:

- $AdS_3$  CS gauge field  $A \leftrightarrow$  chiral fermion current  $J = \psi_L^\dagger \Gamma \psi_L$
- $AdS_2$  gauge field  $a \leftrightarrow$  impurity spin  $S = \chi^\dagger T \chi$
- $AdS_2$  bifundamental scalar  $\Phi \leftrightarrow$  Kondo interaction  $\mathcal{O} = \psi_L^\dagger \chi$

## Bottom-Up Model

$$S = S_G + S_{CS} + S_{AdS_2}$$

$$S_G = \frac{1}{2\kappa_N^2} \int d^3x \sqrt{-g} \left( R + \frac{2}{L^2} \right)$$

$$S_{CS} = -\frac{\mathcal{N}}{4\pi} \int tr \left( A \wedge dA + \frac{2}{3} A \wedge A \wedge A \right)$$

$$S_{AdS_2} = -\mathcal{N} \int d^3x \delta(x) \sqrt{-\gamma} \left[ \frac{1}{4} tr f^2 + |\mathcal{D}\Phi|^2 + V(\Phi^\dagger\Phi) \right]$$

$$\mathcal{D}\Phi \equiv \partial\Phi + iqA\Phi - iq a\Phi, \quad \Phi = \phi e^{i\psi}, \quad f = da$$

**Bottom-up:** Choose  $V(\Phi^\dagger\Phi) = M^2\Phi^\dagger\Phi$

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**Finite temperature:** BTZ black hole

$$ds^2 = \frac{1}{z^2} \left( \frac{dz^2}{h(z)} - h(z) dt^2 + dx^2 \right), \quad h(z) = 1 - \frac{z^2}{z_H^2}$$

## Field Equations

The equations of motion (EOMs) are

$$\epsilon^{n\mu\nu} F_{\mu\nu} = -\frac{4\pi}{N} \delta(x) J^n$$

$$\partial_m (\sqrt{-g} g^{mp} g^{nq} f_{pq}) = -J^n$$

$$\partial_m (\sqrt{-g} g^{mn} \partial_n \phi) = \sqrt{-g} \Delta^m \Delta_m \phi + \frac{1}{2} \sqrt{-g} \frac{\partial V}{\partial \phi}$$

$$\partial_n J^n = 0$$

where we parametrise  $\Phi = \phi e^{i\psi}$ , and define

$$J^n \equiv 2\sqrt{-g} g^{mn} \Delta_m \phi^2, \quad \Delta_m \equiv A_m - a_m + \partial_m \psi$$

## Static Solution

Ansatz ( $a_z = 0$  gauge):  $a_t(z)$ ,  $\phi(z)$ ,  $A_x(z) \neq 0$

Near the boundary  $z \rightarrow 0$ :

$$a_t(z) \sim \frac{Q}{z} + \mu, \quad \phi(z) \sim \sqrt{z} [\alpha \log(\Lambda z) + \beta], \quad A_x(z) \rightarrow 0$$

$|\mathcal{O}|^2$  marginal  $\leftrightarrow M^2 - Q^2 = -1/4 = AdS_2$  BF bound

**Double trace coupling** [Witten 01]:  $\alpha = \kappa\beta \propto \langle \mathcal{O} \rangle$

- $\kappa$  Kondo coupling
- $\kappa$  runs:  $\phi(z)$  independent of  $\Lambda$



## Running of Kondo Coupling

$T = 0$ : Under scale change  $\Lambda \rightarrow \Lambda'$ ,

$$\kappa(\Lambda') = \frac{\kappa(\Lambda)}{1 + \kappa(\Lambda) \log(\Lambda/\Lambda')}$$

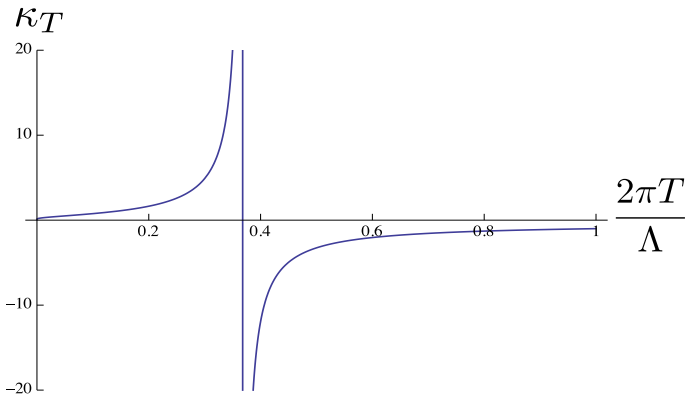
For  $\kappa < 0$ :

- **Asymptotic freedom**:  $\Lambda'/\Lambda \rightarrow \infty$ ,  $\kappa \rightarrow 0$
- **IR divergence**:  $\Lambda'/\Lambda \rightarrow 0$ ,  $\kappa \rightarrow \infty$  at  $\Lambda' = \Lambda e^{1/\kappa(\Lambda)}$

$T \neq 0$ : Under rescaling  $x^\mu \rightarrow \tilde{x}^\mu z_H$ ,  $a_t \rightarrow \tilde{a}_t/z_H$ ,  $z_H = (2\pi T)^{-1}$

$$\Phi(z) \rightarrow \tilde{\Phi}(\tilde{z}) \sim \alpha_T \sqrt{\tilde{z}} \log(\tilde{z}) + \beta_T \sqrt{\tilde{z}}$$
$$\kappa_T \equiv \frac{\alpha_T}{\beta_T} = \frac{\kappa}{1 + \kappa \log(\Lambda z_H)} \quad (\text{RG invar.})$$

## Dynamical Scale Generation

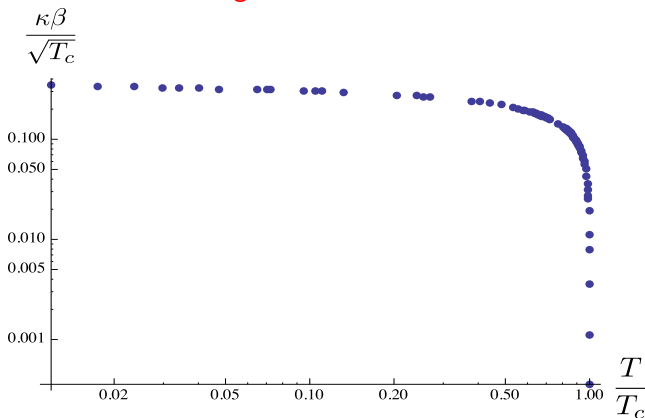


$$\kappa(\Lambda) = -1$$

Divergence  $\Rightarrow T_K = \frac{1}{2\pi} \Lambda e^{1/\kappa(\Lambda)}$ ,  $T/T_K = e^{-1/\kappa_T}$  (**RG invar.**)

# Phase Transition

Large- $N$  Kondo effect



$$\langle \mathcal{O} \rangle \propto \kappa\beta \propto \left(1 - \frac{T}{T_c}\right)^{1/2}, \quad T \lesssim T_c \approx 0.89T_K, \quad \kappa T_c \approx 8.98$$

# Quantum Quenches

Study the response of the holographic system as the Kondo coupling is varied in time

## **Quench protocols:**

- Gaussian pulse
- Constant slope, erf, tanh

## **Consider:**

- Quench within same phase
- Quench across phases

## Eddington-Finkelstein Coordinates

The numerical problem is most conveniently solved in the **Eddington-Finkelstein (EF)** coordinates where perturbations can travel from the boundary to the horizon in finite coordinate time.

Define EF time coordinate by

$$dv = dt - \frac{dz}{h(z)} \Rightarrow v = t - \operatorname{arctanh} z$$

The background metric becomes

$$ds^2 = \frac{1}{z^2} \left( -h(z)dv^2 - 2dvdz + dx^2 \right)$$

## EF Coordinate EOMs

Ansatz:  $A_x(z, v)$ ,  $a_t(z, v)$ ,  $a_z(z, v)$ ,  $\phi(z, v)$ ,  $\psi(z, v) \neq 0$

$$0 = \partial_v A_x - 4\pi\delta(x)\phi^2 (\partial_v \psi - h\partial_z \psi)$$

$$0 = 4\pi\delta(x)\phi^2 \left( \partial_z \psi - \frac{a_v}{h} \right) + \partial_z A_x$$

$$0 = -\frac{\partial_v^2 a_v}{h} + \partial_z \partial_v a_v - \frac{2\phi^2 (\partial_v \psi - h\partial_z \psi)}{z^2}$$

$$0 = \frac{(zh' - 2h) \partial_v a_v}{zh^2} - \frac{\partial_z \partial_v a_v}{h} + \frac{2\partial_z a_v}{z}$$

$$+ \partial_z^2 a_v - \frac{2a_v \phi^2}{z^2 h} + \frac{2\partial_z \psi \phi^2}{z^2}$$

$$0 = \phi \left( -\frac{2a_v \partial_v \psi}{h^2} + \frac{a_v^2}{h^2} + \frac{2\partial_v \psi \partial_z \psi}{h} - \partial_z \psi^2 \right)$$

$$+ \frac{h' \partial_z \phi}{h} - \frac{2\partial_z \partial_v \phi}{h} + \partial_z^2 \phi$$

# Solving the PDEs

## General strategy:

- 1 **Solve the initial static problem**
  - Find the initial static solution to the EOMs to the specified accuracy using both shooting and pseudospectral methods.
- 2 **Time-march**
  - Time-evolve the initial solution using the implicit Crank-Nicholson method.

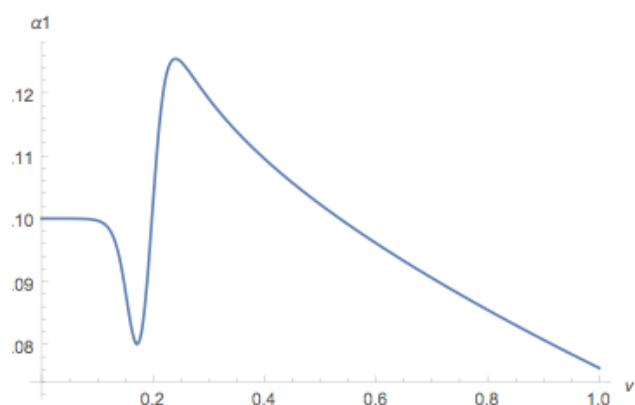
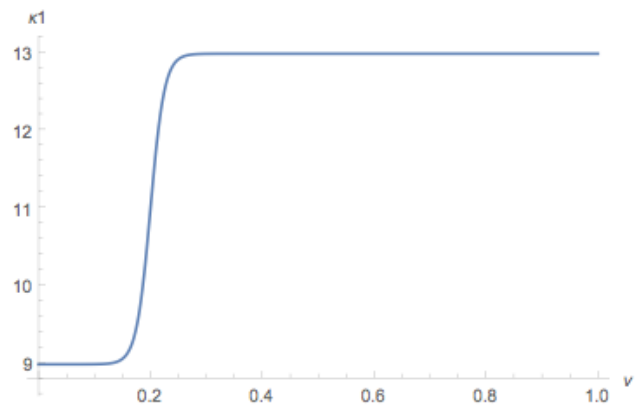
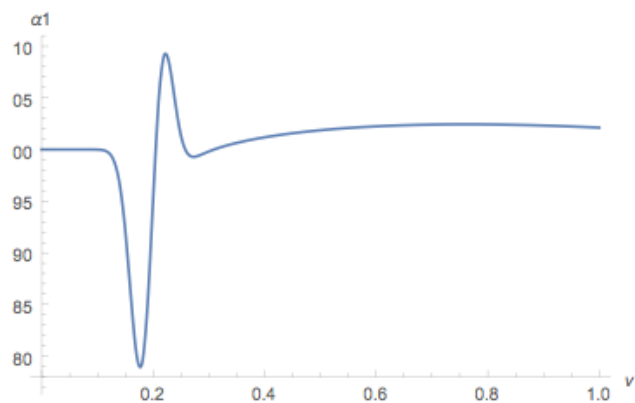
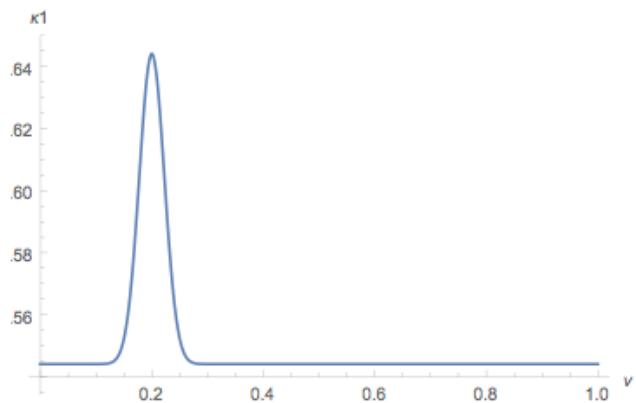
## Need to regulate boundary non-analyticities:

- Use change of variables and field redefinition to remove  $\sqrt{z}$  and  $(\log)^n$  terms so that fields are regular at the boundary up to the second derivatives.

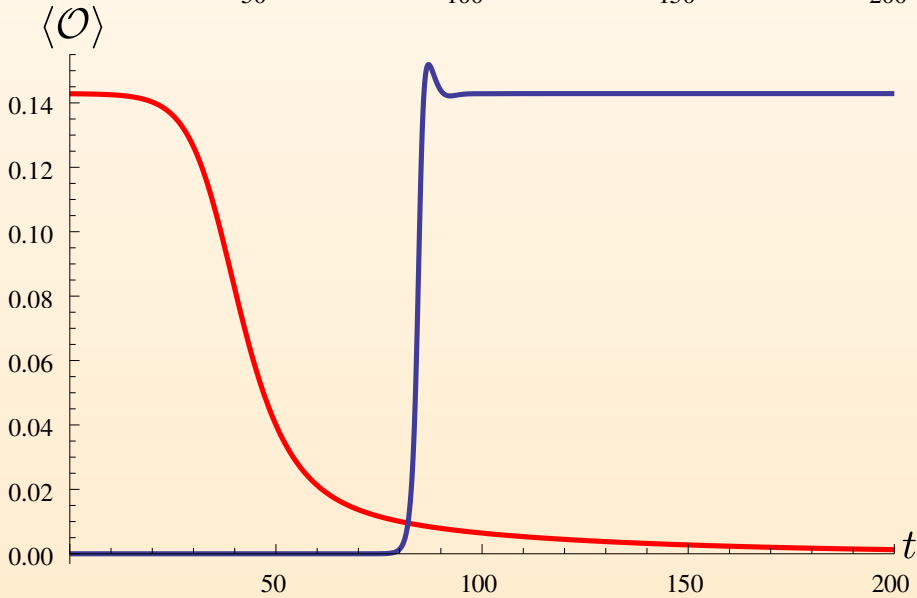
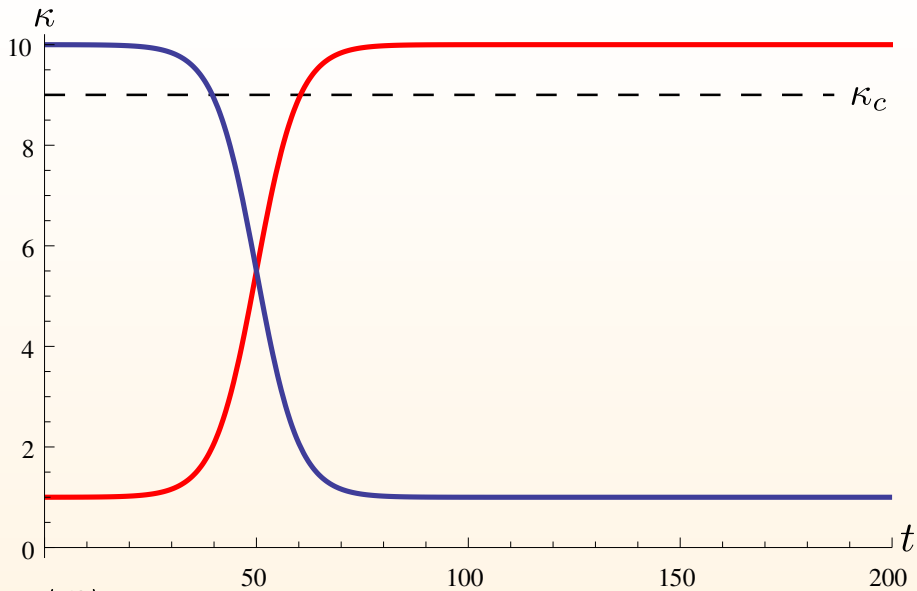
# Kondo coupling



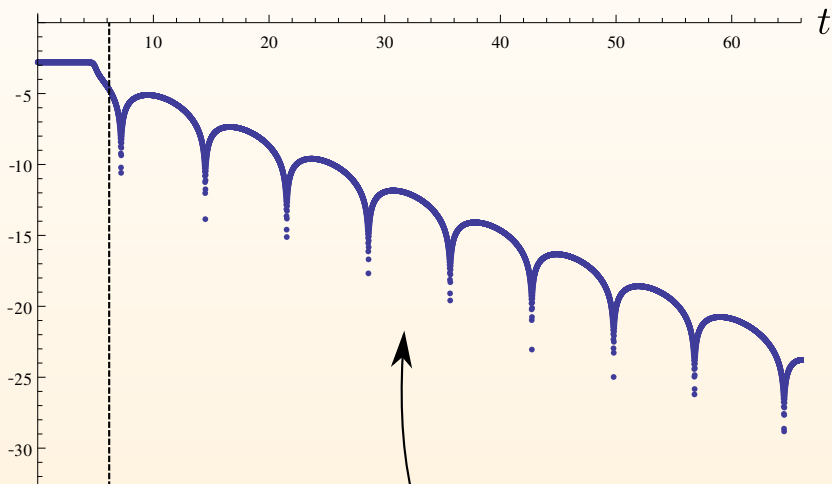
# Condensate



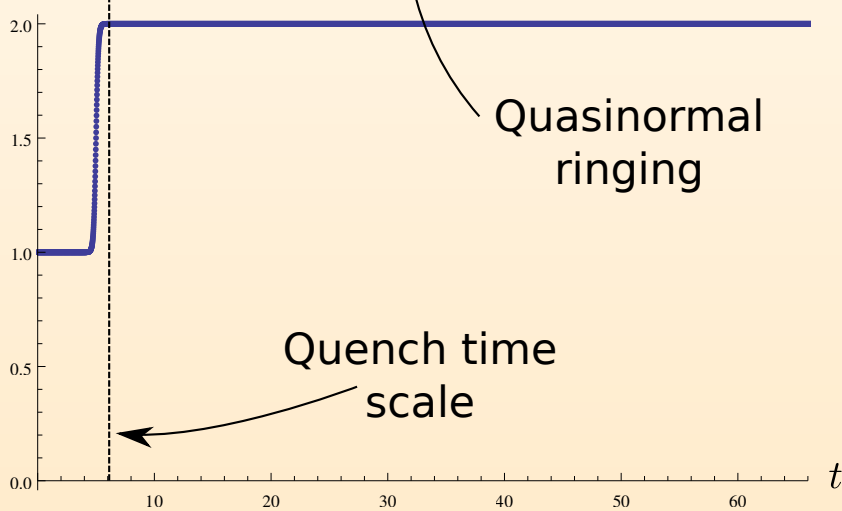




$\log(|\langle \mathcal{O} \rangle|)$



$\kappa(t)$



Quasinormal  
ringing

Quench time  
scale

## Discussion

- Holographic Kondo system relaxes back to the equilibrium state characterised by the end value of the Kondo coupling.
- Relaxation time scale depends on the final size of the condensate.
- Quasi-normal modes determine the relaxation behaviour and control the time scales of the equilibration process.

Cf. [Bayat, Bose, Johannesson, Sodano, PRB 92, 155141 \(2015\)](#): Late-time behaviour of a two-impurity Kondo spin-chain model after a quantum quench characterised by single-frequency oscillations.