



Holographic Heavy Ion Collisions in Non-Conformal Theories

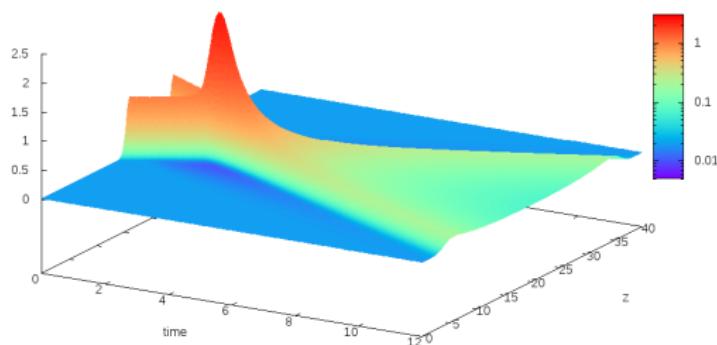
Maximilian Attems

[arXiv:1603.01254](https://arxiv.org/abs/1603.01254)

[arXiv:1604.06439](https://arxiv.org/abs/1604.06439)

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Ioannis Papadimitriou, Daniel Santos, Carlos Sopuerta,
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NumHol2016

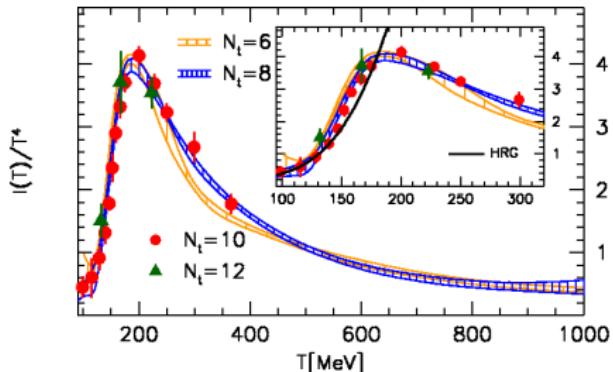


Far from equilibrium dynamics:

- at strong coupling
- fast hydrodynamization time
- IC hydrodynamics

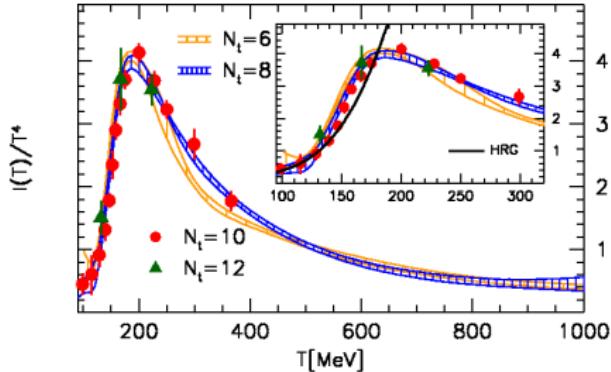
Non-conformal holographic shockwaves

Motivations II

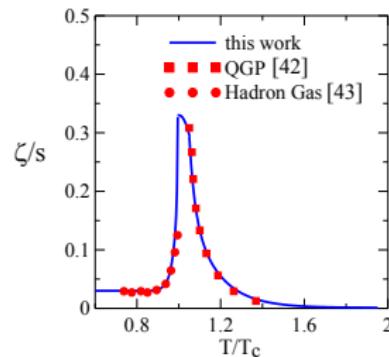


[S. Borsanyi *et alii* arXiv:1007.2580 [hep-lat]]

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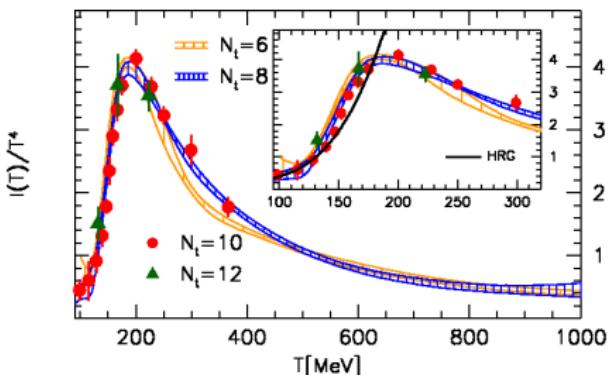


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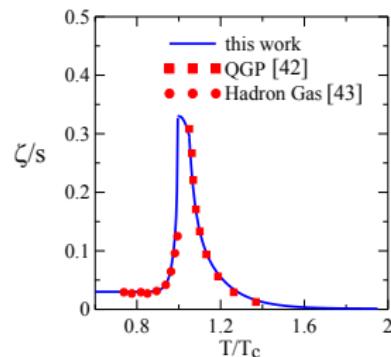


[Denicol *et alii* arXiv:0903.3595 [hep-ph]]

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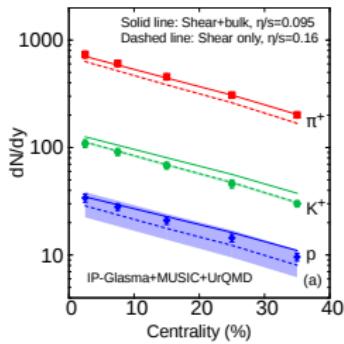


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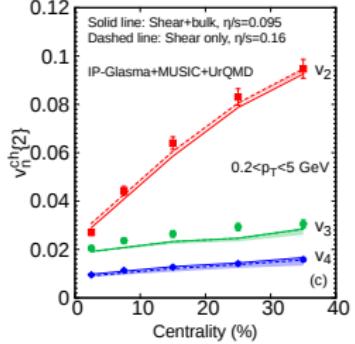
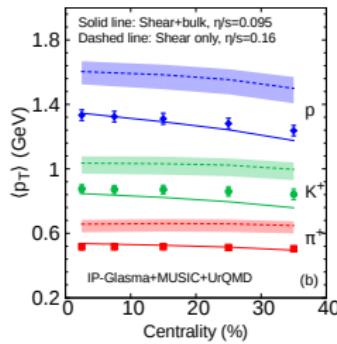


[Denicol *et alii* arXiv:0903.3595 [hep-ph]]

Hydro simulation agreement improves with bulk viscosity:



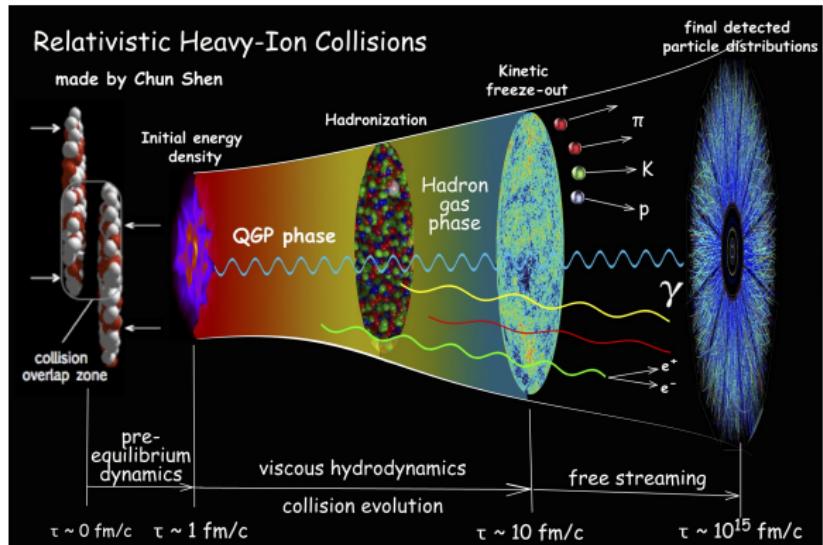
[Denicol *et alii* arXiv:1502.01675 [nucl-th]]



1 Non-conformal shockwave collisions

- Heavy-Ion collision
- General Relativity setup
- Scalar potential
- Interaction measure
- Bulk viscosity
- Buchel bound
- Quasi-Normal-Modes
- Shockwaves Initial Conditions
- Hydrodynamization / equilibration time
- Non-conformal temperature scan

Heavy-Ion collision - little bang



Stages:

- 1) Early out of equilibrium
- 2) Quark-Gluon Plasma
- 3) Particularization

Can we describe all the stages at strong coupling?
Yes! (up to the last one)

Shockwave collisions in $N=4$ SYM

[Chesler, Yaffe 11; Albacete, Kovchegov, Taliotis 08; Grumiller, Romatschke 08]

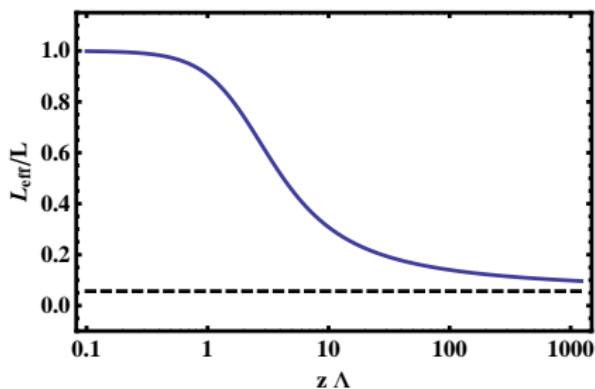
General Relativity setup

Einstein-Hilbert action with scalar potential in five-dimensional bottom-up model:

$$S = \frac{2}{\kappa_5^2} \int d^5x \sqrt{-g} \left[\frac{1}{4} \mathcal{R} - \frac{1}{2} (\nabla\phi)^2 - V(\phi) \right].$$

Potential $V(\phi)$ interpolating between two AdS with L radius of the UV AdS solution:

$$ds^2 = \frac{L_{\text{eff}}(z)^2}{z^2} (-dt^2 + d\mathbf{x}^2 + dz^2).$$



$V(\phi)$ depends on single parameter ϕ_M , setting non-conformality for this bottom-up model:

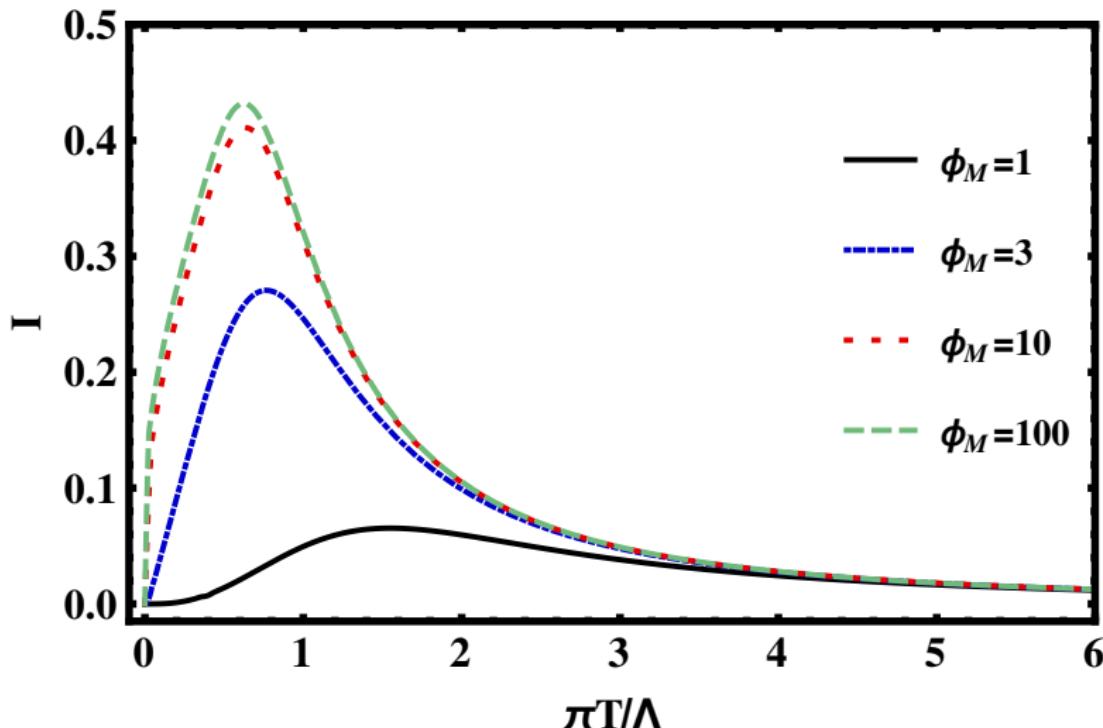
$$L^2 V(\phi) = -\frac{1}{12\phi_M^4} \phi^8 + \left(\frac{1}{2\phi_M^4} + \frac{1}{3\phi_M^2} \right) \phi^6 - \frac{1}{3} \phi^3 - \frac{3}{2} \phi^2 - 3.$$

Deforming $\mathcal{N} = 4$ Super Yang-Mills with an dimension 3 operator \mathcal{O} dual to the scalar field ϕ .

$$\langle T_\mu^\mu \rangle = -\Lambda \mathcal{O}.$$

The source Λ triggers RG flow responsible for the breaking of conformal invariance.

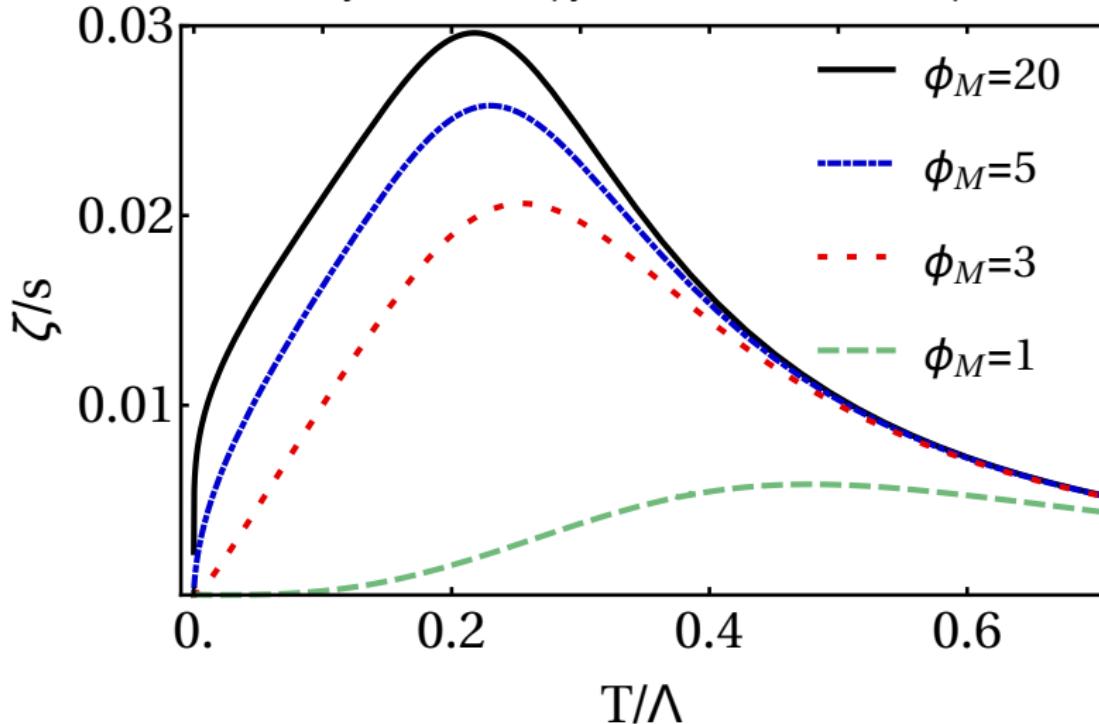
Non-conformality measure as function of the temperature:



conformal at low and high T , **non-conformal** in between

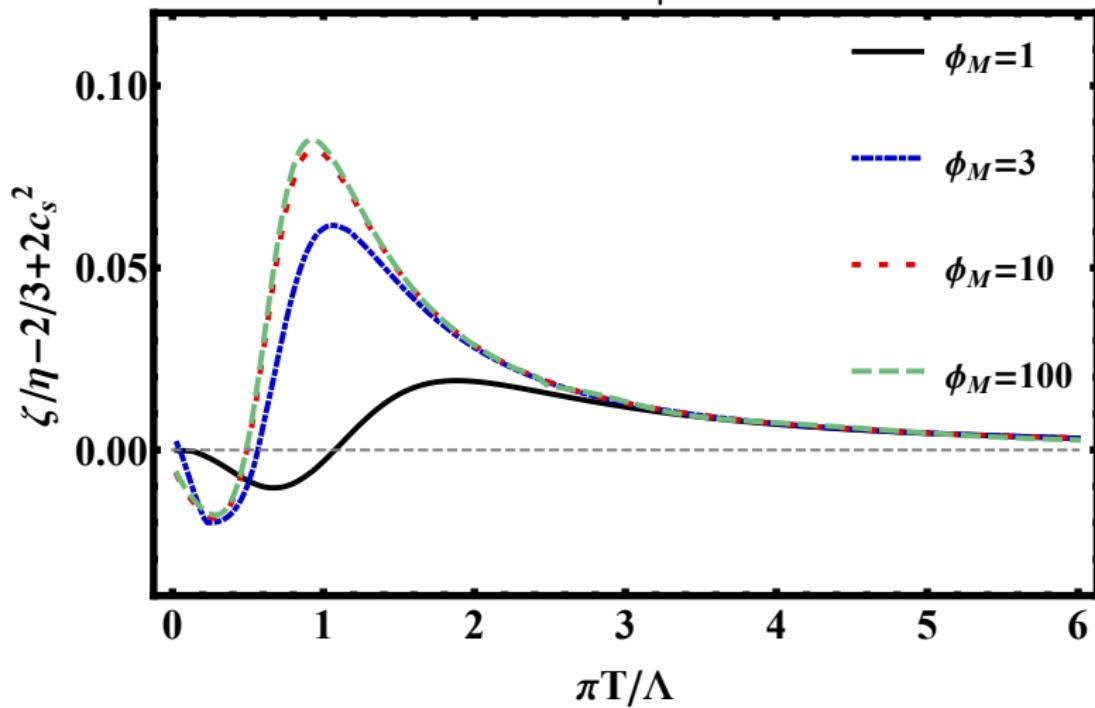
Bulk Viscosity

Ratio of bulk viscosity over entropy as a function of temperature:



non-conformal behaviour reflects in transport coefficients

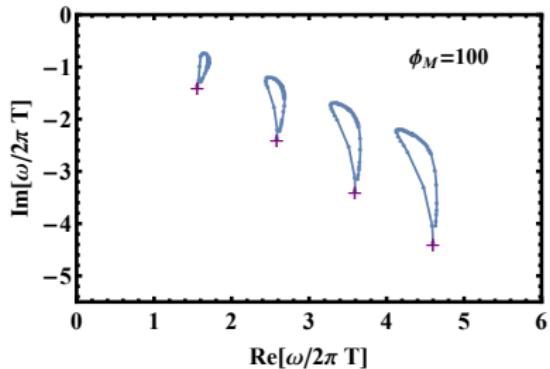
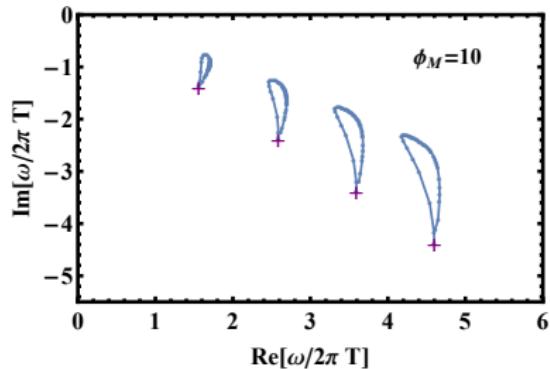
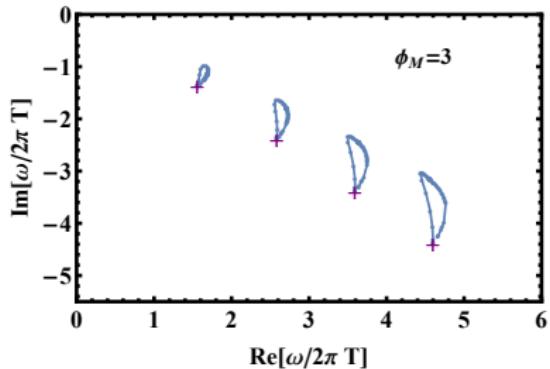
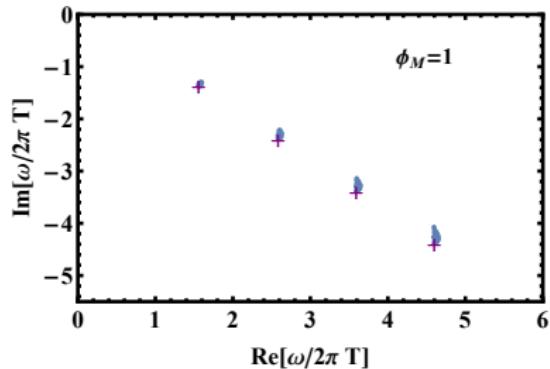
Violation of Buchel's bound at low temperatures:



Maxima of speed of sound and bulk to shear viscosity different!

Quasi-Normal-Modes

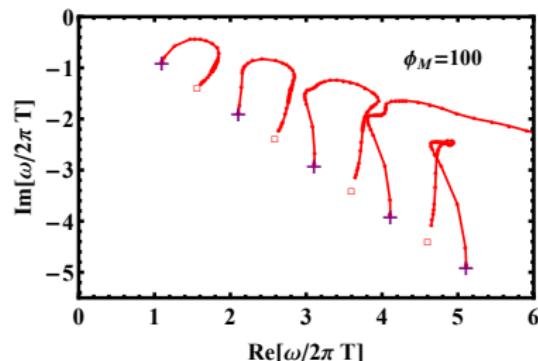
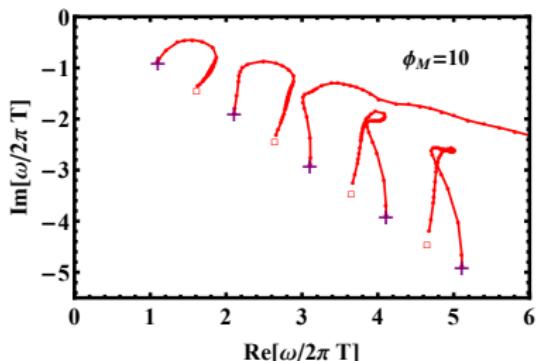
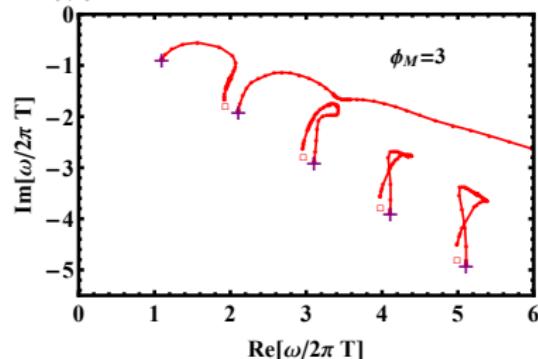
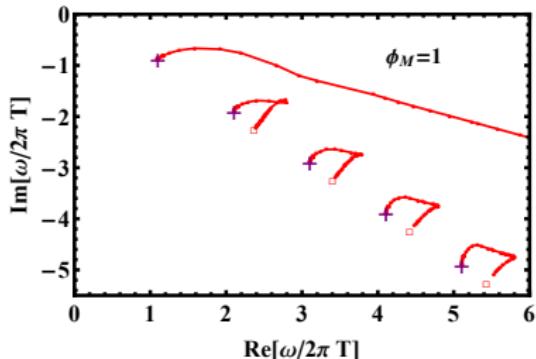
anisotropic perturbation $Z_{\text{aniso}} = e^{-2A}(h_{zz} - h_{aa})$



Fluctuations of the stress energy tensor

Quasi-Normal-Modes

non-conformal mode $Z_{\text{bulk}} = \phi - \frac{e^{-2A(\phi)}}{2A'(\phi)} h_{aa}$



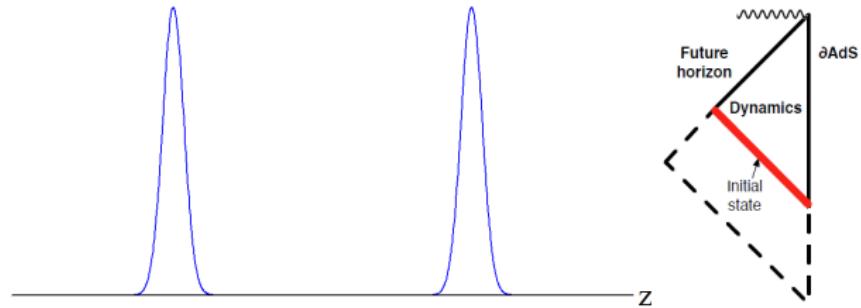
n -th scalar mode decoupling with anti-crossing

Shockwaves Initial Conditions

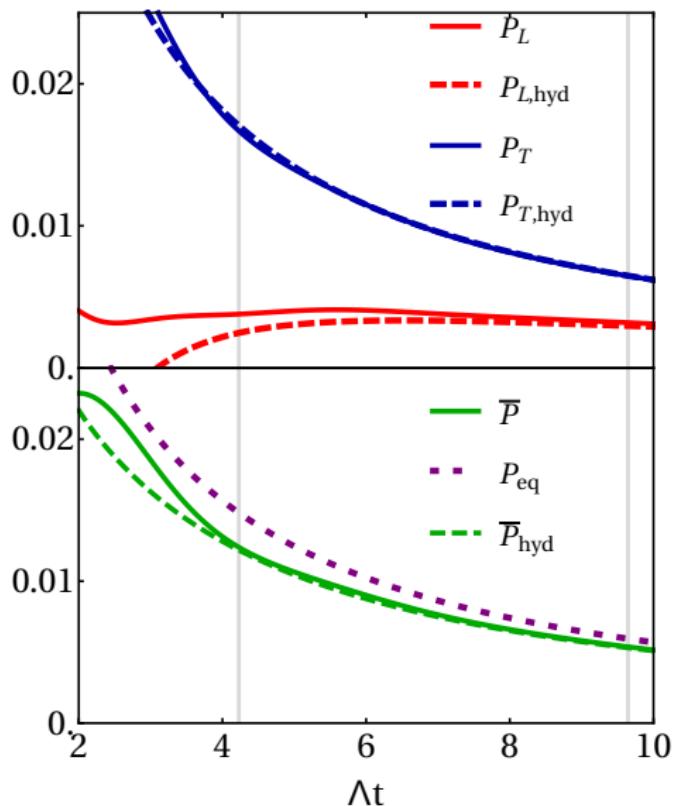
5D metric Ansatz in Eddington-Finkelstein:

$$ds^2 = -A dt^2 + \Sigma^2 \left(e^B dx_\perp^2 + e^{-2B} dz^2 \right) + 2dt(dr + F dz)$$

- Field theory interpretation:
 - Defined by energy density
 - Move in AdS₅ space
 - Demand that shockwaves move at speed of light
 - Quantum state/AdS geometry completely fixed for pure gravity
- Homogeneous in transverse plane ('infinite nucleus')



Hydrodynamization / equilibration time



Hydrodynamics assumes
mean free path goes to zero:

$$\partial_\mu T^{\mu\nu} = 0$$

$$T^{\mu\nu} = (\epsilon + p)u^\mu u^\nu + pg^{\mu\nu} + \eta\Pi^{\mu\nu} + \zeta\Pi(g^{\mu\nu} + u^\mu u^\nu)$$

Hydrodynamization:

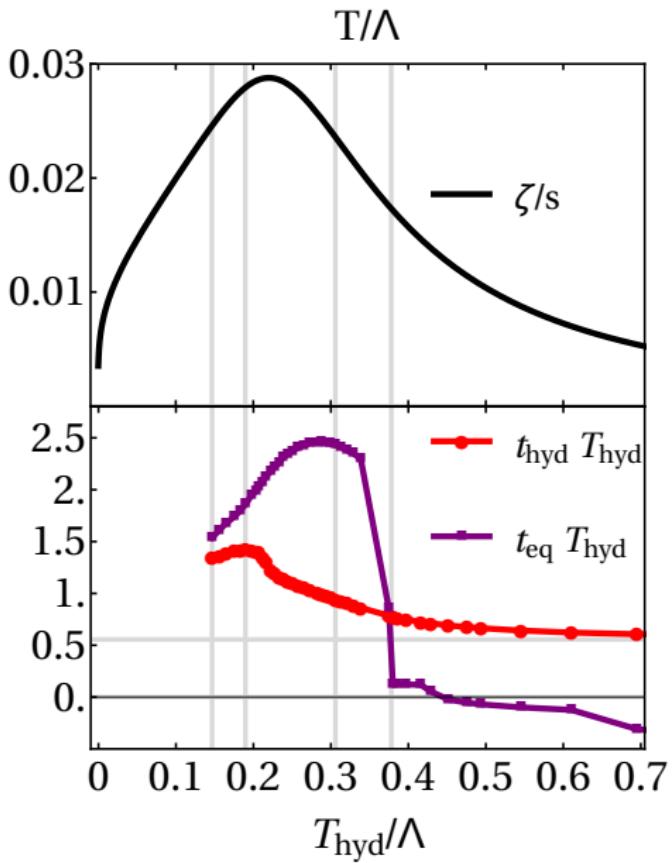
$$\left| P_{L,T} - P_{L,T}^{\text{hyd}} \right| / \bar{P} < 0.1$$

Equilibration:

$$\left| \bar{P} - P_{eq} \right| / \bar{P} < 0.1$$

hydrodynamization \neq equilibration \neq isotropization

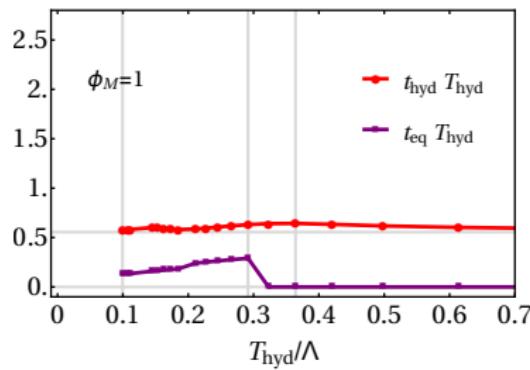
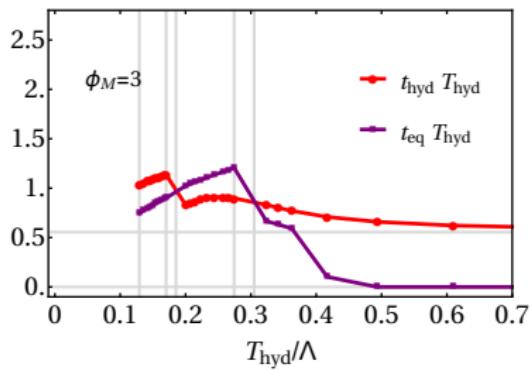
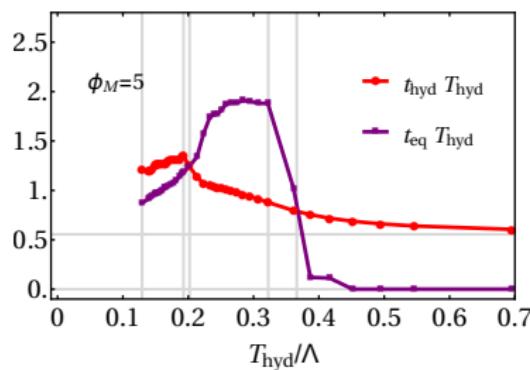
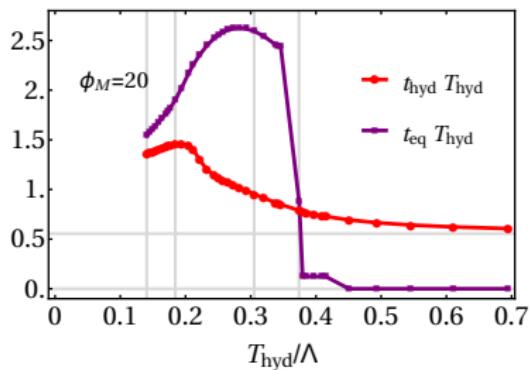
Non-conformal temperature scan



Non-conformal T scan:

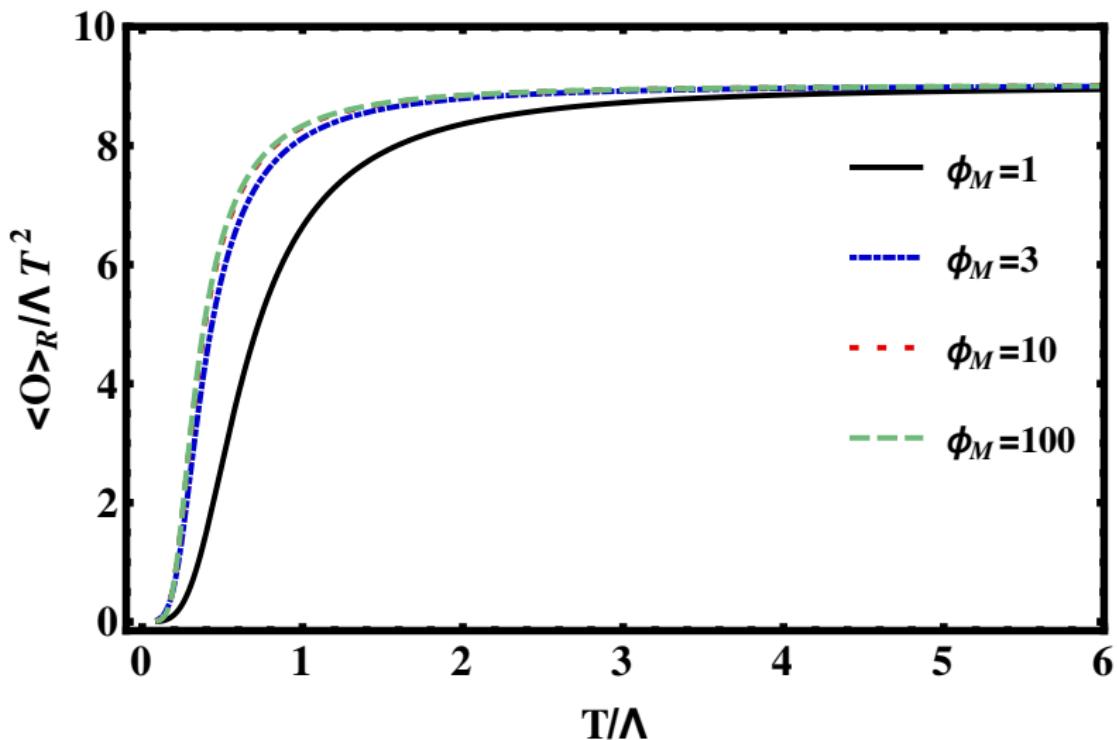
- EOS does NOT hold out of equilibrium
- t_{hyd} slow down
- ordering of t_{eq} and t_{hyd} depends on bulk viscosity
- required ζ 1/10 of QCD at T_c

Non-conformal temperature scan



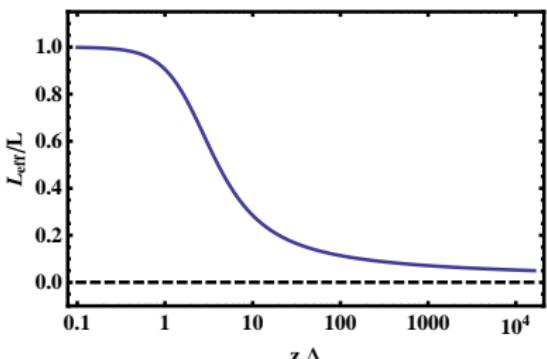
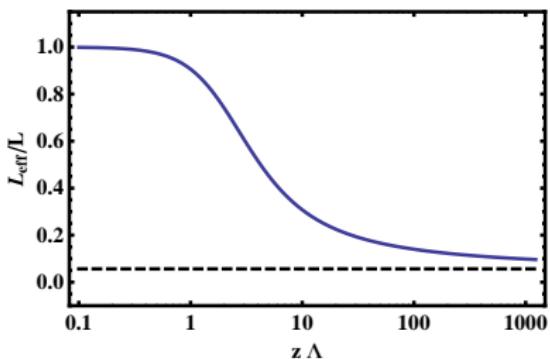
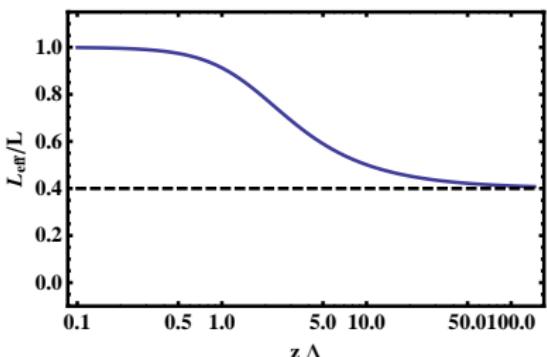
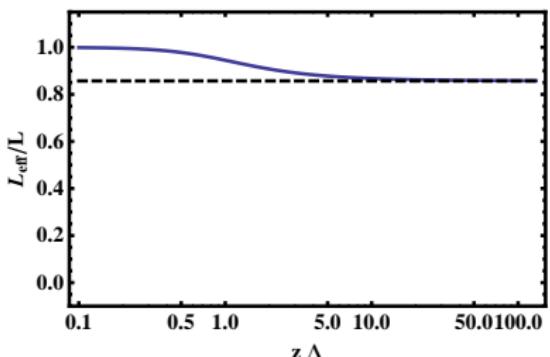
$$\zeta/s \approx 0.02 \text{ for } t_{\text{eq}} > t_{\text{hyd}}$$

- First simulation of a holographic **non-conformal** model for heavy ion collisions
- New relaxation channel from bulk viscosity:
hydrodynamization without equilibration
- Hydrodynamics works at early time (though max delayed ≈ 3)
 - despite non-trivial equation of state
 - despite sizeable ζ/s bulk viscosity over entropy
- More studies are on the way:
 - Systematic exploration of parameter space on *MareNostrum*
 - Asymmetrical collisions
 - Holographic energy scan
 - Different potentials



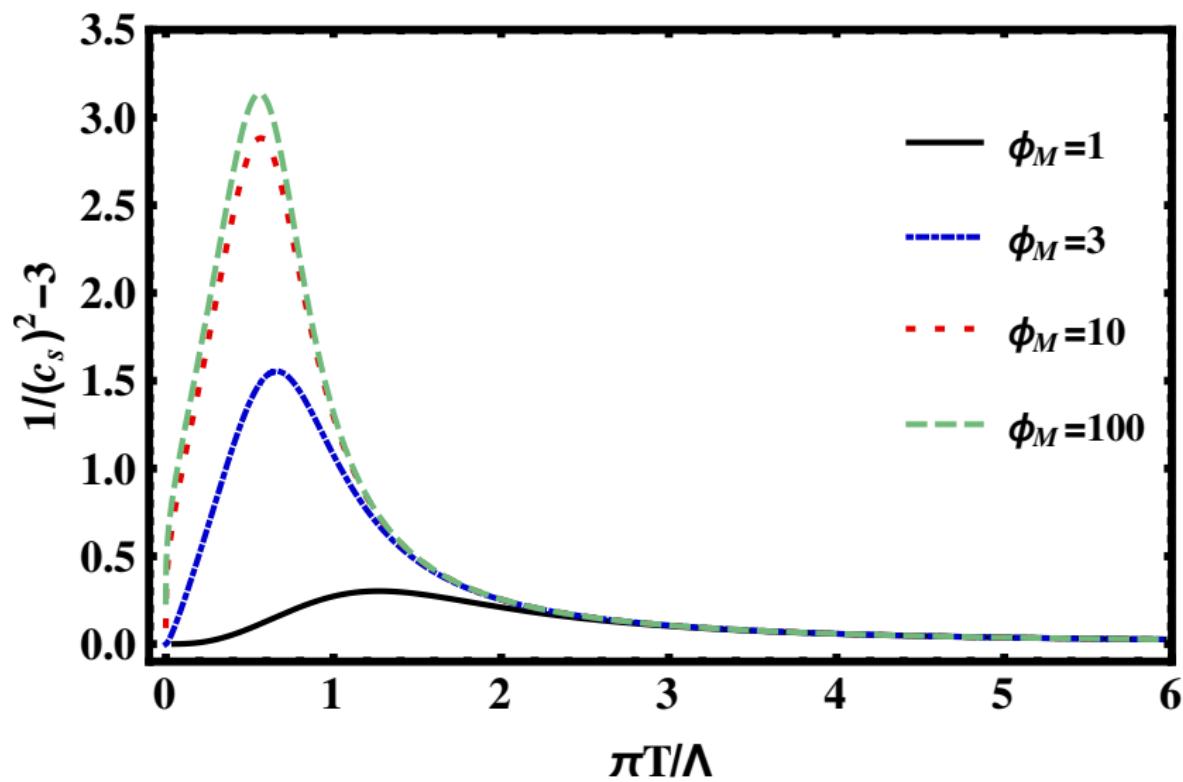
Temperature dependence of the VEV of the scalar operator $\langle \mathcal{O} \rangle_T$ for several values of ϕ_M . $\langle \mathcal{O} \rangle_R = \kappa_5^2 \langle \mathcal{O} \rangle_T / L^3$, $\epsilon - 3p = \Lambda \langle \mathcal{O} \rangle_T$.

Backup: Non-conformal effects



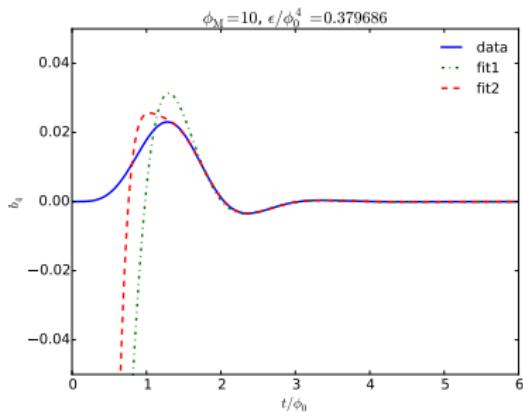
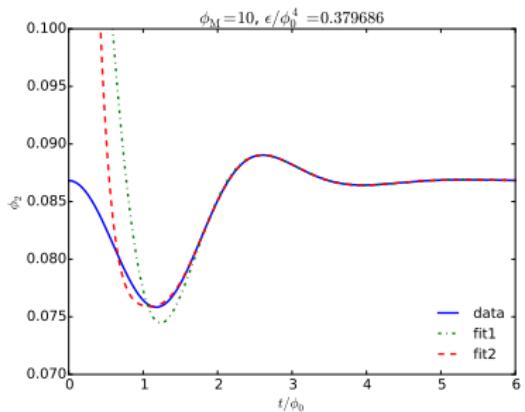
$$ds^2 = \frac{L_{\text{eff}}(z)^2}{z^2} (-dt^2 + d\mathbf{x}^2 + dz^2) .$$

Backup: Speed of sound



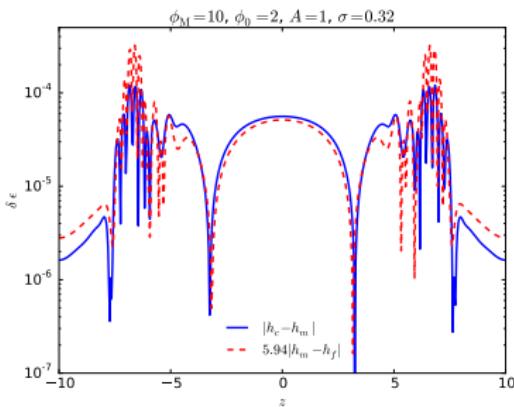
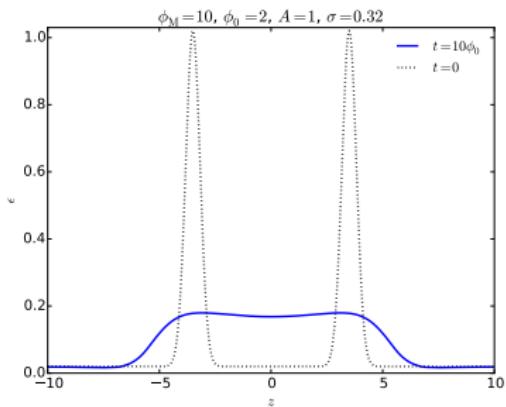
Inverse speed of sound square as a function of T

ϕ_2 and b_4 as functions of time for a z -independent configuration



Blue full line corresponds to data from the code, green dash-dotted line correspond to a fit to the data using one QNM, red dashed line corresponds to a fit using two QNMs as explained in the text.

Differences between the coarse and medium (blue solid line) and the medium and fine (red dashed line) resolution run



The results show fourth-order convergence.