Probing new physics in $B_s \rightarrow \mu^+ \mu^-$ at LHCb

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Outline

- \Box Motivation of $B_s \rightarrow \mu \mu$ as a probe of NP
- □ Analysis in LHCb
 - Overview
 - Calibration and normalization
- First results
- Conclusions



Motivation of $B_s\!\!\rightarrow\!\!\mu\mu$ as a probe of NP



- \Box $B_s \rightarrow \mu\mu$ can access NP through new virtual particles entering in the loop \rightarrow indirect search of NP
- Indirect approach can access higher energy scales and see NP effects earlier:
 - Done before in the history of physics...
 - 3rd quark family inferred by Kobayashi and Maskawa (1973) to explain CPV in K mixing (1964). Directly observed in 1977 (b) and 1995 (t)
 - Neutral Currents discovered in 1973, Z⁰ directly observed in 1983



Decay Physics (SM)

Hadronic weak decays are often studied in terms of effective hamiltonians of local operators. Degrees of freedom of exchanged particles are integrated out giving rise to the Wilson coefficients C_i.

SUSY...)

$$BR(B_{q} \to \mu^{+} \mu^{-}) = \frac{G_{F}^{2} \alpha^{2}}{64\pi^{3}} |V_{tb}^{*} V_{tq}|^{2} \tau_{Bq} M_{Bq}^{3} f_{Bq}^{2} \sqrt{1 - \frac{4m_{\mu}^{2}}{M_{Bq}^{2}}} \times \left\{ M_{Bq}^{2} \left(1 - \frac{4m_{\mu}^{2}}{M_{Bq}^{2}}\right) C_{s}^{2} + \left[M_{Bq} C_{P} + \frac{2m_{\mu}}{M_{Bq}} C_{10} \right]^{2} \right\}$$

- C_{S,P} are negligible in SM, C₁₀ gives the only relevant contribution.
 - This decay is very suppressed in SM:

 $\mathbf{C}_{\mathbf{P},\mathbf{S},\mathbf{10}}$ (pseudoscalar, scalar and axial) depend on the underlying model (SM,



 $BR(B_s \to \mu \mu) \ = \ (3.35 \pm 0.32) \cdot 10^{-9} \ \text{M.Blanke et al., JHEP 10 003, 2006}$

Current experimental upper limit (CDF, 2fb⁻¹) still one order of magnitude above these values. @ 90% CL:

BR(B_s \rightarrow µµ) < 3.6·10⁻⁸ CDF collab., CDF Public Note 9892

* BR (Bd $\rightarrow \mu\mu$) = (1.03 ± 0.09)·10⁻¹⁰

New Physics effects

- □ NP can contribute to this decay rate (specially SUSY at high tan β (tan β = v_u/v_d)):
- $\hfill\square$ More than one Higgs \rightarrow contributions to $\boldsymbol{C}_{\boldsymbol{S},\boldsymbol{P}}$
 - 2HDM-II : BR proportional to tan⁴β
 - SUSY (MSSM): above + extra tan⁶β +...
- RPV SUSY: tree level diagrams
- Technicolor (TC2), Little Higgs (LHT) ... modify C₁₀.



NP can modify the BR from smaller SM up to current experimental upper limit \rightarrow **Any measured value will affect NP searches!**



Analysis in LHCb

 \rightarrow Overview



LHCb overview





VELO

RICH-1

Magnet

LHCb overview

RICH-2

Tracker

OT

LHCb strong points:

M1

ECal

- PID
- Vertexing and IP
- Momentum and mass resolution

Calorimeters

HCal

- Flexible trigger

SPD

Physics: Examples of key channels

- Rare decays: $B_s \rightarrow \mu \mu$, $B_d \rightarrow K^* \mu \mu$

PS

- CP Violation: $B_s \rightarrow J/\Psi \Phi$, $B_{s/d} \rightarrow hh$

Muon System



VELO

RICH-1

LHCb overview

RICH-2

Tracker

TT

OT

More on LHCb, see: - LHCb status and physics plans for 1fb⁻¹ run by Eugeni Grauges - Flavor tagging at LHCb by Marc Grabalosa

Magnet

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Muon System



- □ Selection: apply some cuts on all µµ candidates to remove most of the background.
- Classify each event using three properties (bins in a 3D parameter space):
 - Particle Identification (PID):
 Probability to be muons
 - Geometrical properties (Geometrical likelihood)
 - Invariant Mass
- 3D space is binned, so that each bin is treated as an independent experiment. Results combined using Modified Frequentist Approach



- Use of control channels to avoid dependence on simulation:
 - Calibration of relevant variables
 - Normalization

see T. Junk NIM A434, 435,1999

□ How is the Geometrical Likelihood built?

- 1. Input variables: min Impact Parameter Significance (μ^+,μ^-) , DOCA, Impact Parameter of B, lifetime, iso μ^+ , iso- μ^-
 - What is isolation? Fake B_s→µµ might originate in muons from another SV. For each muon, remove the other µ and look at the rest of the event: How many good - SVs (forward, DOCA, pointing) can it make?

Distributions of some relevant variables from MC signal and background







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- 3. In such space correlations are more linear-like \rightarrow rotation matrix, and repeat 2





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- 4. Transformations under signal hyp. $\rightarrow \chi^2_{S}$, under bkg. $\rightarrow \chi^2_{B}$.
- 5. Discriminating variable is $\chi^2_s \chi^2_B$, flat again





PID Likelihood

- Particles with associated hits after extrapolation to the muon chambers are flagged as muons
- Some of them might not be actual muons (misidentification). Different subdetectors return probabilities for different kinds of particles, as seen before:
- Muon chambers: distances of hits to track extrapolation, or fit of the track to hits
- RICH: uses masses of the particles
- CALOs: energy deposition
- Probabilities can be combined in a likelihood to fight against remaining misid





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Sensitivity (I)

□ Signal 90% CL **exclusion sensitivity** as a function of Luminosity and time



Sensitivity (II)

□ Signal **evidence sensitivity** as a function of Luminosity and time





Analysis in LHCb

\longrightarrow Calibration and normalization



- □ Signal is distributed in several bins of a 3D space
- We need to know both the overall normalization and the fraction of signal in each bin
 - **Invariant mass:** Can be calibrated, e.g., with fit of $B \rightarrow hh$ line shape or from charmonium and bottonium resonances
 - **Geometrical Likelihood**: $B \rightarrow hh$ triggered independently of signal (event triggered by the other B)
 - **PID likelihood:** J/ Ψ taking p, pt distributions from B \rightarrow hh
- Will have a quick look at invariant mass and geometrical likelihood

Invariant mass calibration (I)

□ Method 1: Full fit of the B → hh line shape

– Use the Geometrical Likelihood to clean the sample!



		GL:[0.25-0.5]	GL:[0.5-0.75]	GL:[0.75-1.0]
	Ntot	703	661	695
	σ (MeV/c²)	24.4 ± 2.6	27.5±0.82	26.9±0.2
	$M(B_d)$ (MeV/c ²)	5276.6±4.3	5273±4.0	5273±0.71
	$M(B_s)$ (MeV/c ²)	5352.8±3.8	5359±3.2	5356±0.77

Invariant mass calibration (II)

Method 2: Interpolation of σ between charmoniums and bottomiums.





Geometrical Likelihood calibration (I)

 \Box Fit the distribution in GL bins to extract the **number of B** \rightarrow hh signal events.

- First bin difficult to fit. Get the number in this bin by:
 - Fitting all $B \rightarrow hh$ and subtracting events from other GL bins,
 - Cleaning with topological cuts before the fit and assessing signal loss via control channel.

 $\sigma = 38.0 \pm 2.2 \text{ MeV/c}^2$

 $N\pi\pi_{sig} = 640 \pm 37$

 $N_{bkg} = 2670 \pm 78$

 N_{bka} phys = 84 ± 45

GL=[0.5-

5800

0.751

5600

5400



Events / (57.5 MeV/c²

500

400

300

200

100

4800

LHCb

Preliminary

s = 7 TeV Data

5000

5200

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 \Box GL distribution from B \rightarrow hh:



Normalised into a pdf, used in the computation of the limit



□ Normalization needed to convert # events into a BR without relying on knowledge of σ_{bb} , integrated luminosity or absolute efficiencies

 $BR = BR_{cal} \cdot \frac{\varepsilon_{cal}^{\text{Rec}} \cdot \varepsilon_{cal}^{\text{Sel}} \cdot \varepsilon_{cal}^{\text{GEC}} \cdot \varepsilon_{cal}^{\text{Trig}}}{\varepsilon_{Bs}^{\text{Rec}} \cdot \varepsilon_{Bs}^{\text{Sel}} \cdot \varepsilon_{Bs}^{\text{Sel}} \cdot \varepsilon_{Bs}^{\text{GEC}} \varepsilon_{Bs}^{\text{Trig}}} \cdot \frac{f_{cal}}{f_{B_s}} \cdot \frac{N_{B \to \mu\mu}}{N_{cal}} = norm \cdot N_{B \to \mu\mu} \qquad cal: \text{ control channel} \\ B_s: B_s \to \mu\mu$

Using any, B⁺, B^d as a control channel implies a ~13 % systematic from the knowledge of f_d/f_s (=f₊/f_s). Normalization to a B_s mode would introduce in principle larger errors because of worse known of branching ratios.

□ Some control channel candidates:

- B⁺→J/Ψ(µµ)K⁺, B_s→J/Ψ(µµ)Φ(KK)
 - similar trigger and PID,
 - different reconstruction because of the extra track/tracks
 - B_s : worse BR precision, but not f_+/f_s
- B→hh:
 - Same kinematics
 - different trigger & PID

Several groups in LHCb are measuring f_d/f_s . Hope to reduce the error soon to $\sim 7\%$

Normalization example: $B^+ \rightarrow J/\Psi(\mu\mu)K^+$

- The fraction of efficiencies (trigger, reconstruction, selection, PID...) needs to be computed/cancelled.
- □ Trigger on **data**:
 - J/ Ψ trigger efficiency L0,HLT1 and HLT2



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Normalization example: $B^+ \rightarrow J/\Psi(\mu\mu)K^+$

- The fraction of efficiencies (trigger, reconstruction, selection, PID...) needs to be computed/cancelled.
- □ Reconstruction/acceptance:
 - Ratio 4/3 body allows to evaluate with data the ratio 3/2 body
 - Extra track effect: (acceptance + tracking)

$$\frac{\varepsilon_{2body}^{rec}}{\varepsilon_{3body}^{rec}} \approx \frac{\varepsilon_{3body}^{rec}}{\varepsilon_{4body}^{rec}} = 0.58 \pm 0.04 \quad (stat \ only)$$





First results

□ For each bin we need to know:

 Number of *expected signal events* (for a given BR & Lumi) and the number of *expected background events* (for a given Lumi)
 The number of *observed* events from data (blinded until January!)

For expected limit

we generate toy experiments from signal and background expectations.





□ No $B_s \rightarrow \mu\mu$ data from blinded region used. But all the pdfs (signal and background **mass** and **GL**) obtained from **data**, no MC.





Conclusions



Conclusions

 \square BR(B_s \rightarrow µµ) can constraint **several NP models.**

- Value allowed from current experimental upper limit to below SM prediction.
- Analysis in LHCb based in a 3D (Geometrical Likelihood, Invariant Mass and Particle Identification) space.
- Calibration of Invariant mass and Geometrical Likelihood needed to determine the fraction of signal in each bin. Normalization to control channels used to calculate BR.
- LHCb working really fine. Data approaching MC in momentum and mass resolution, IP, and PID.
- Experimental limit by LHCb close to current's world best at the end of this year. Hope to overtake it soon!



Backup





[arXiv:0708.2079v4 [hep-ph]] (2008)

Maximal CP Violating Minimal Flavour Violation: Enhancements up to current upper limit, but also < SM depending on the phases

CMSSM: departures from SM possible, but less likely taking into account current experimental constraints

New Physics effects (II)



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